

# Reflectance Models (BRDF)

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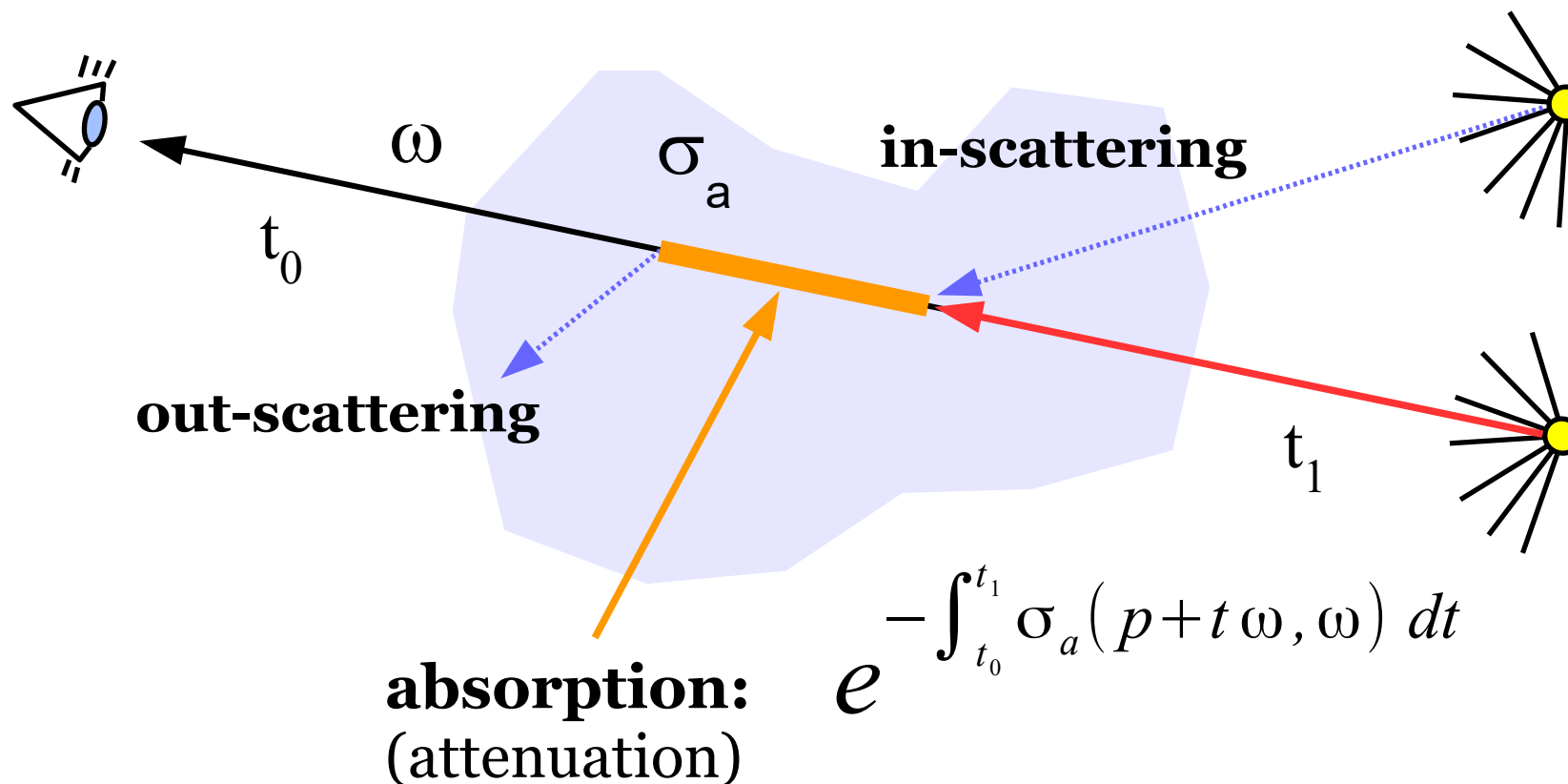
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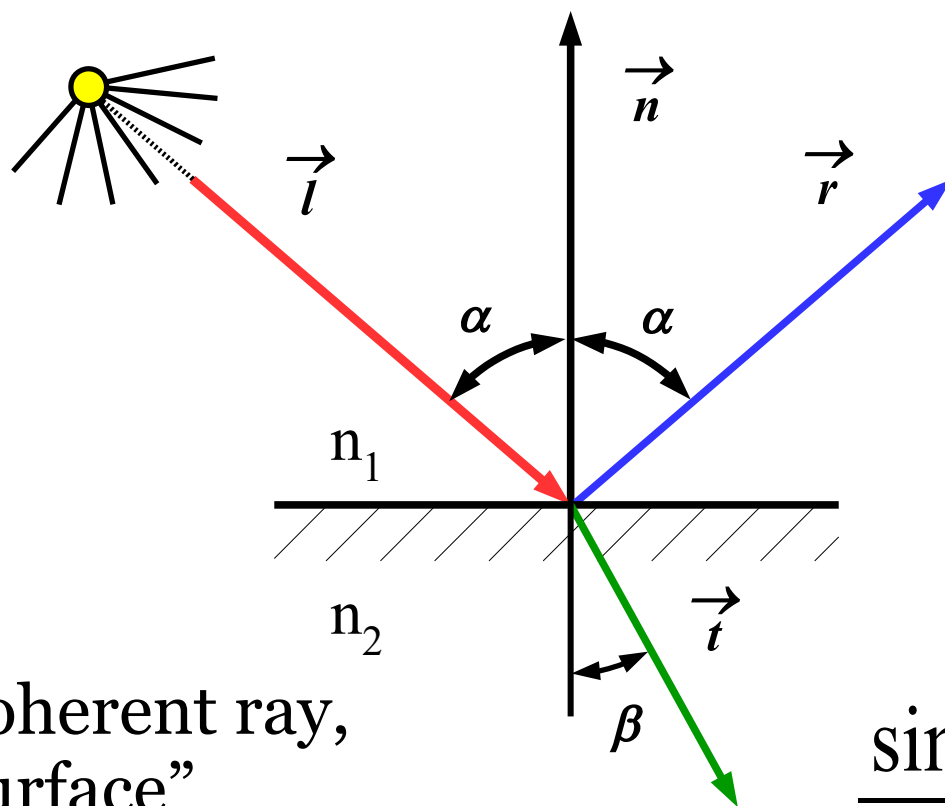
# Light travels through media



Absorption is simple, ~~scattering~~ is very complicated



# Light hits object surface (ideal)



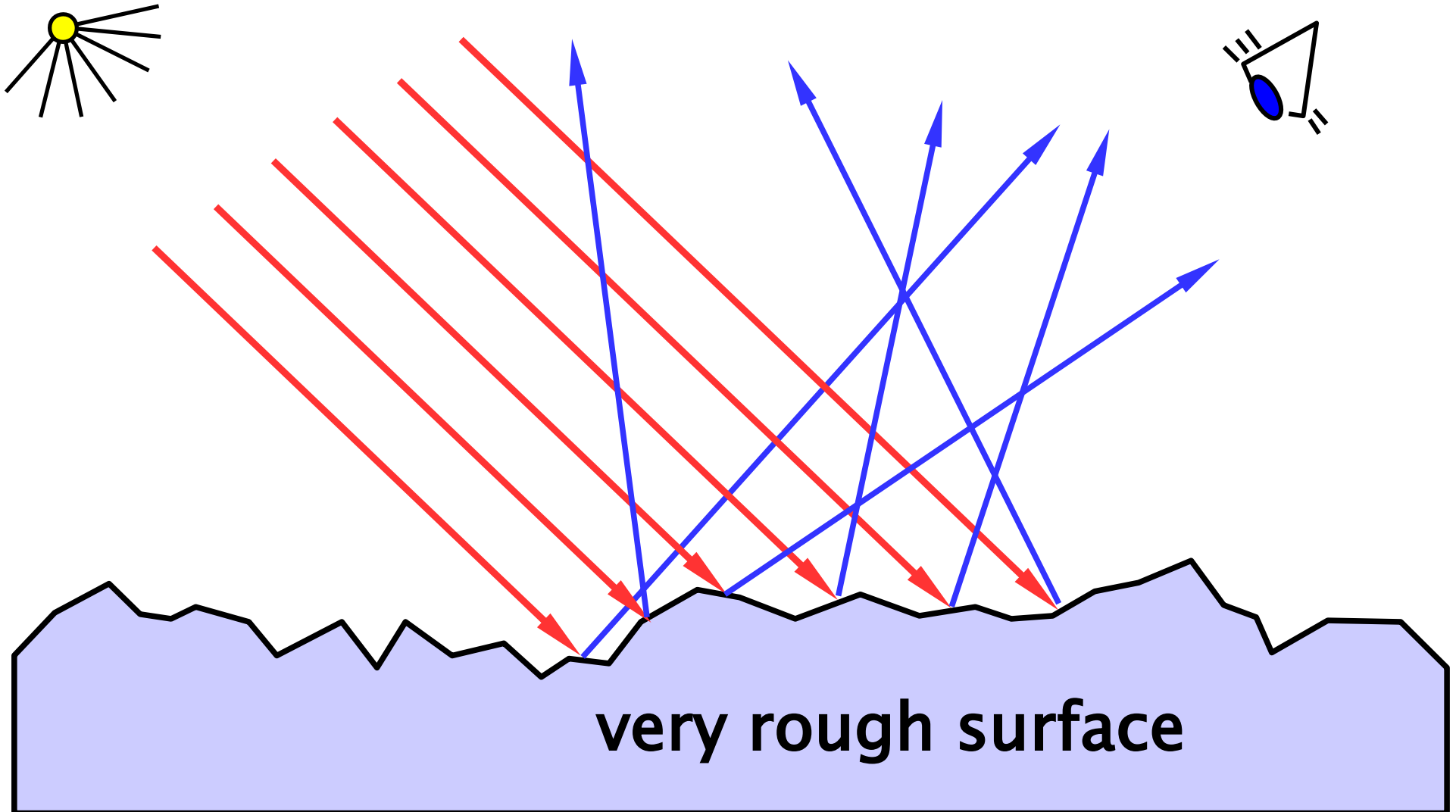
$$r = 2n(n \cdot l) - l$$

“Ideal coherent ray,  
ideal surface”

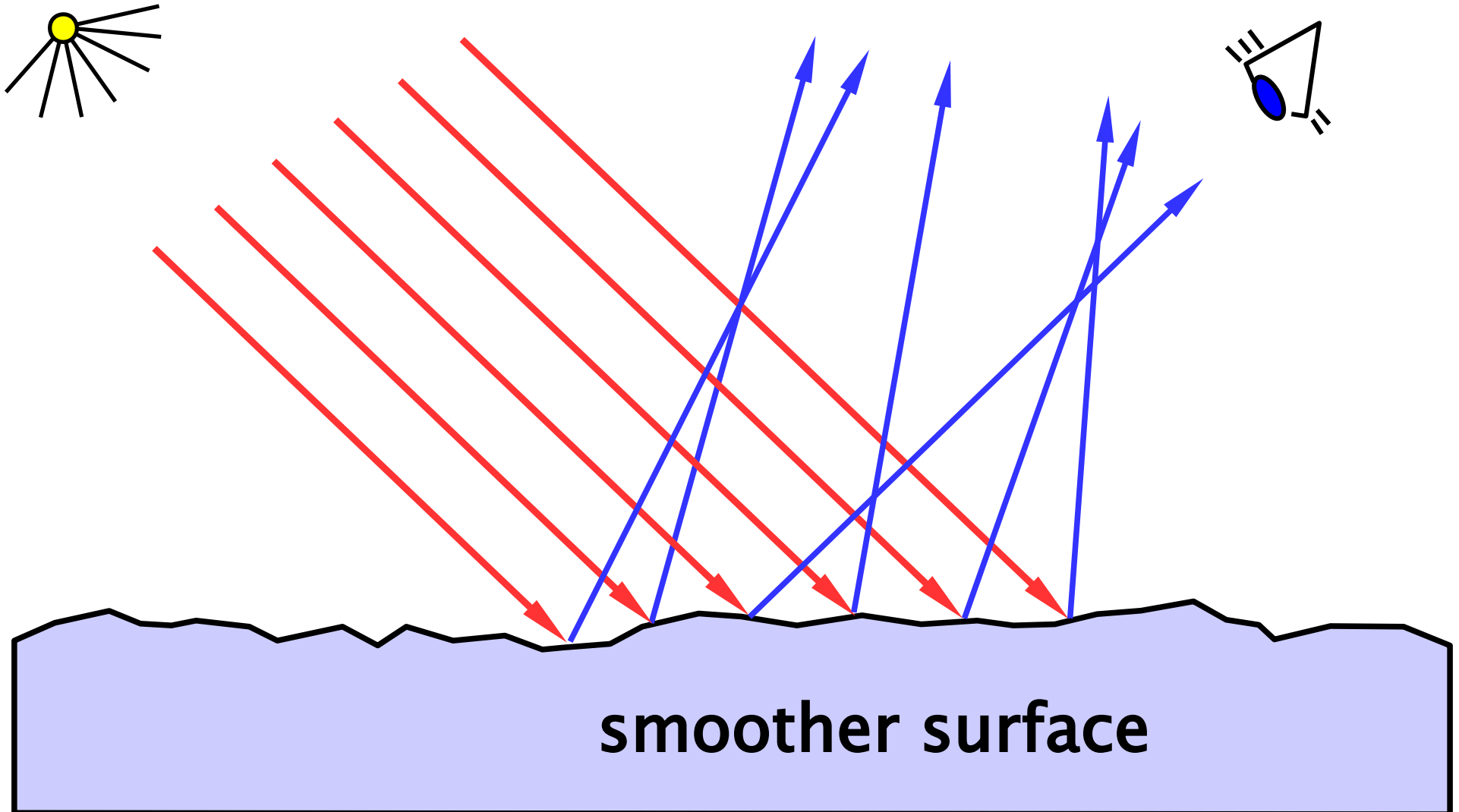
$$\frac{\sin \beta}{\sin \alpha} = \frac{n_1}{n_2} = n_{12}$$

## Fresnel equations

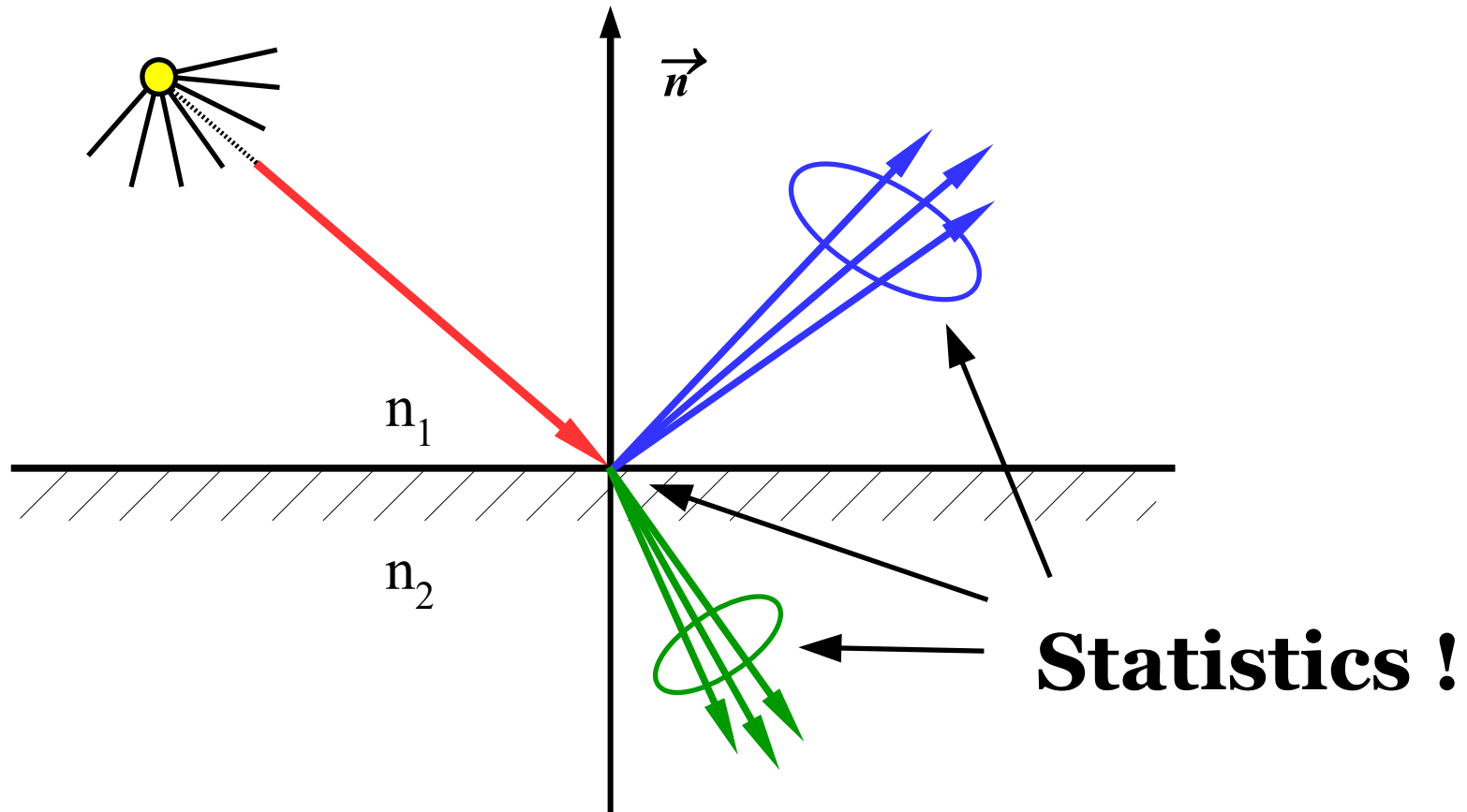
# Real surface (microscopic view)



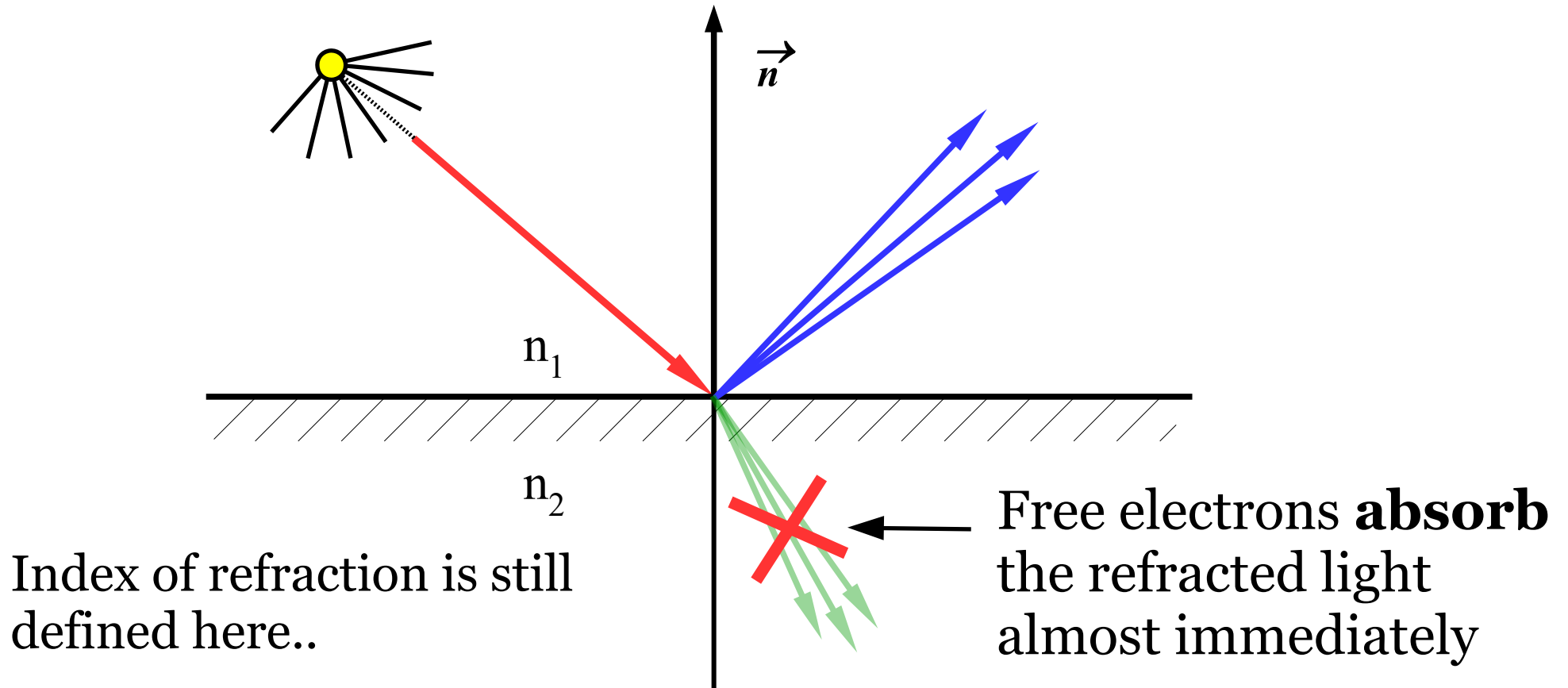
# Real surface (microscopic view)



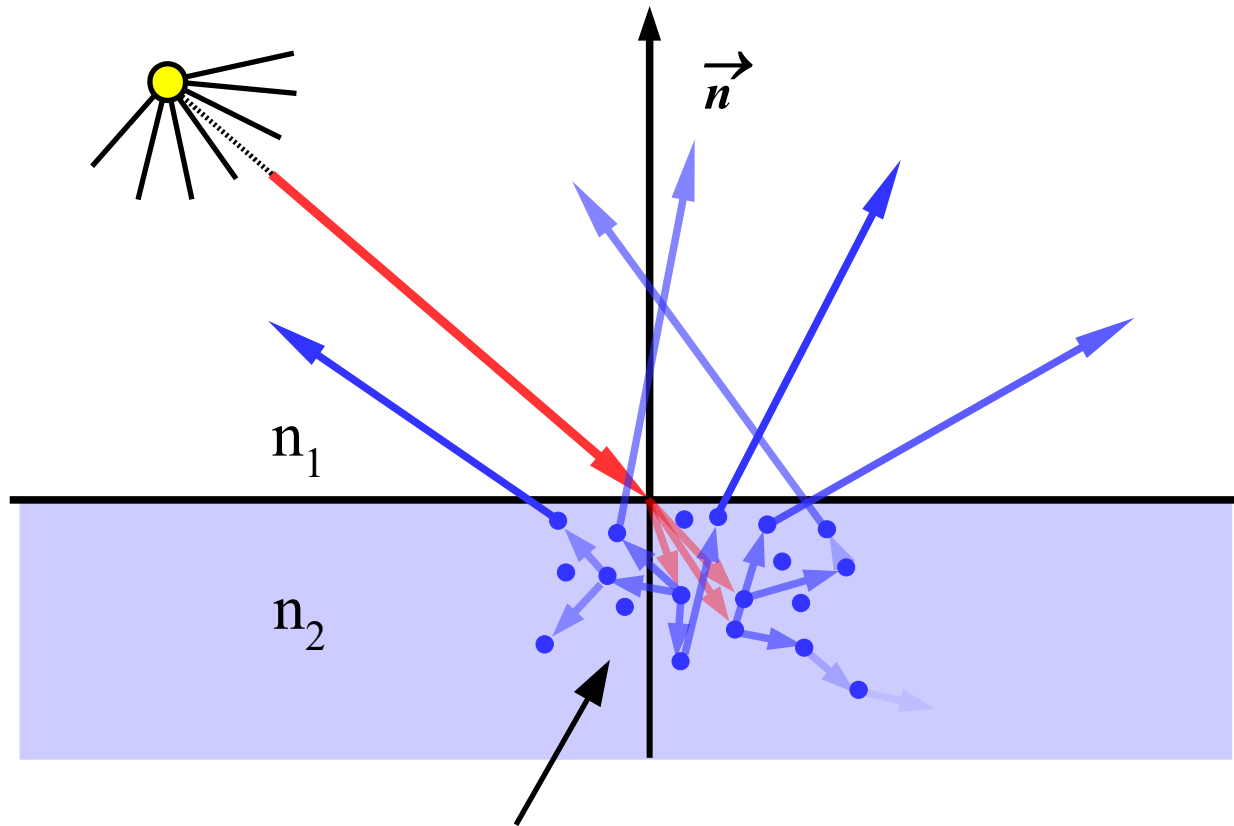
# What we see from the distance..



# Metals (conductors)



# Dielectrics (insulators)

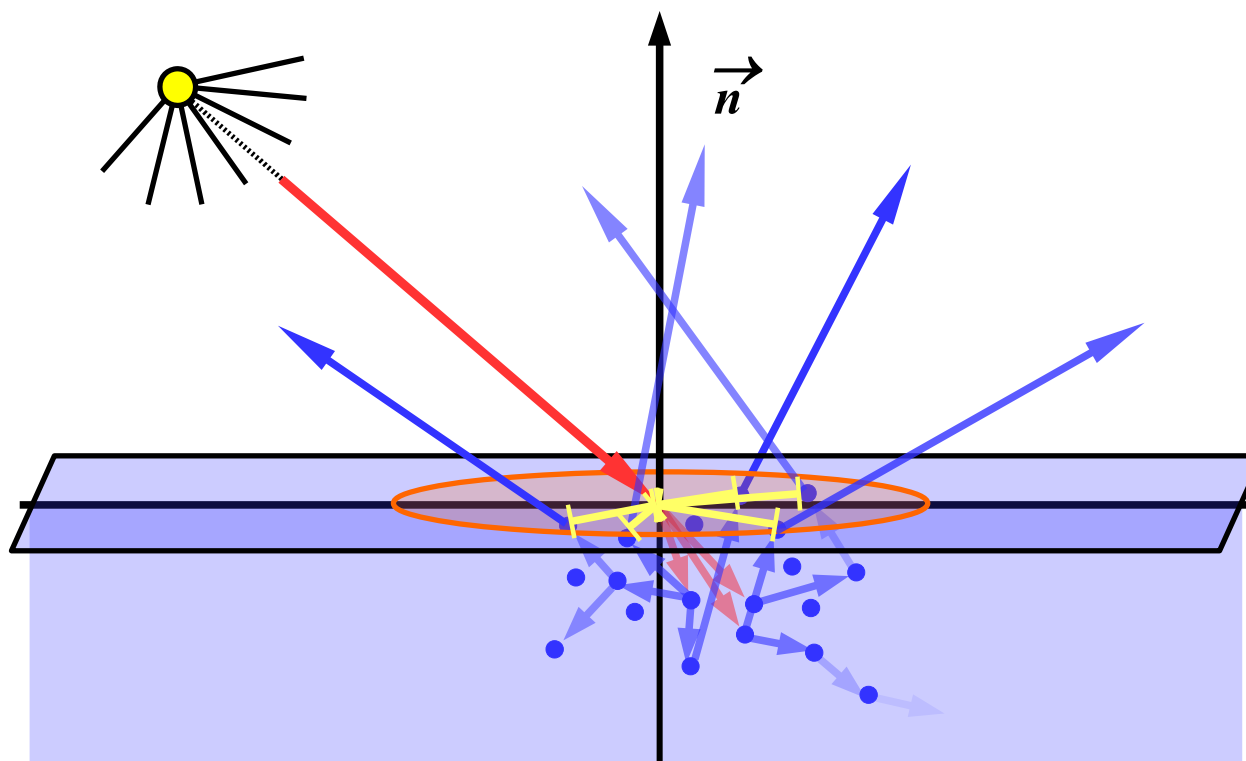


**Scattering** inside of the material



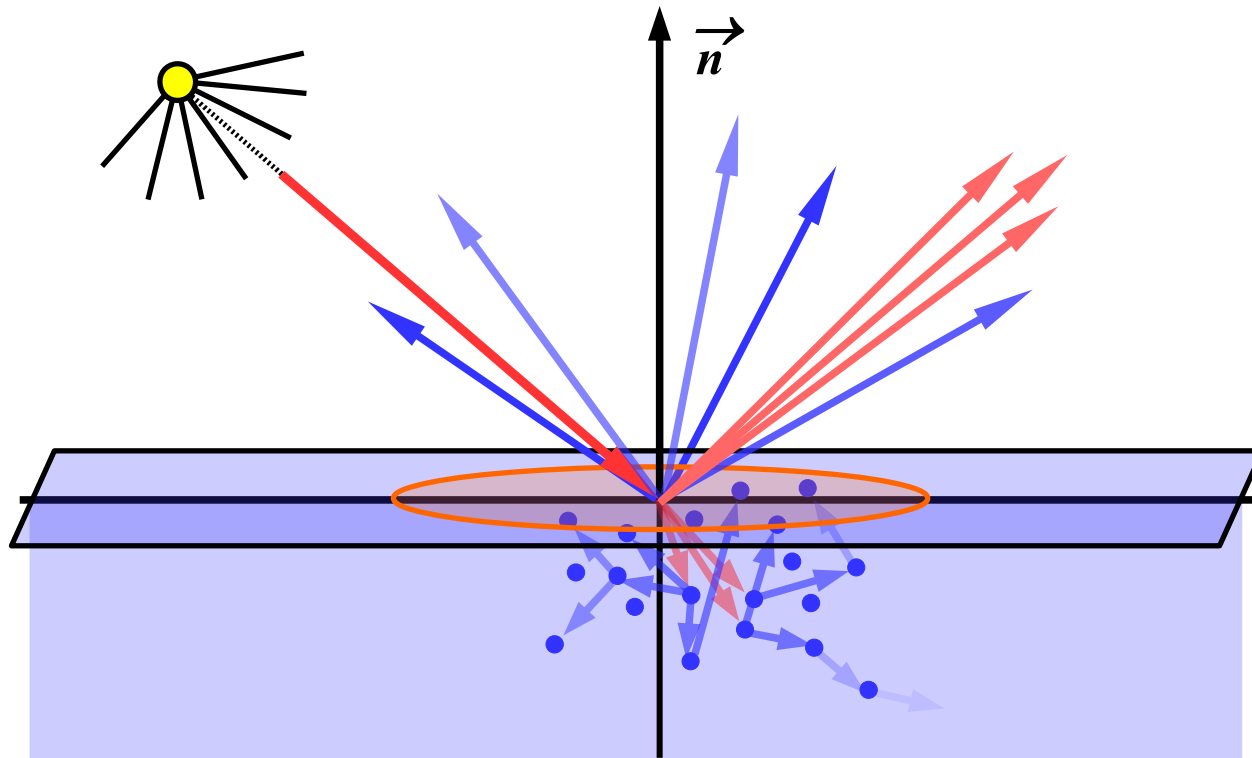


# BSSRDF idea



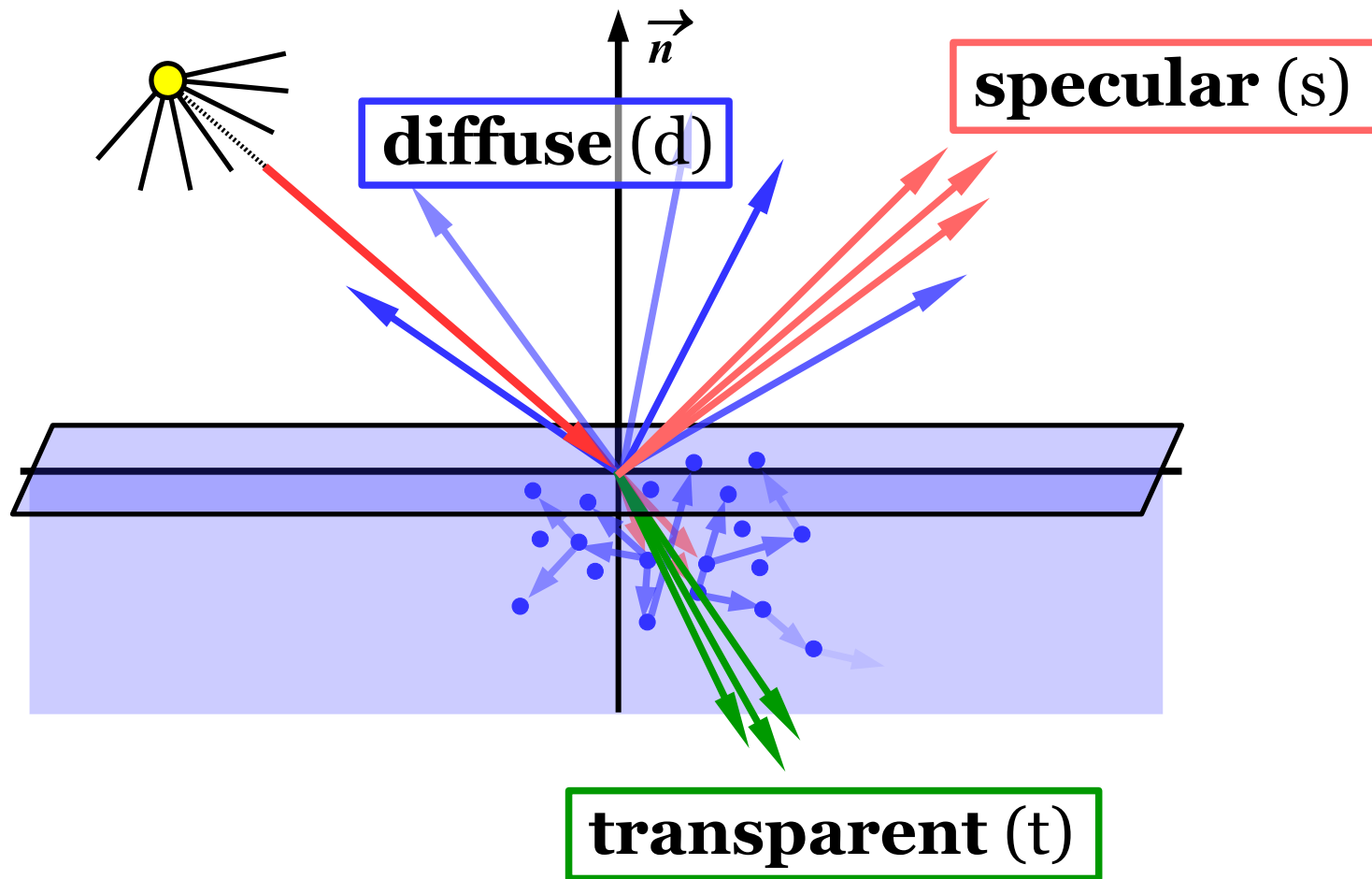
(“Bi-directional Scattering-Surface Reflectance Distribution Function”)

# Ignoring exit-to-entry distance



**BRDF** = “Bi-directional Reflectance Distribution Function”

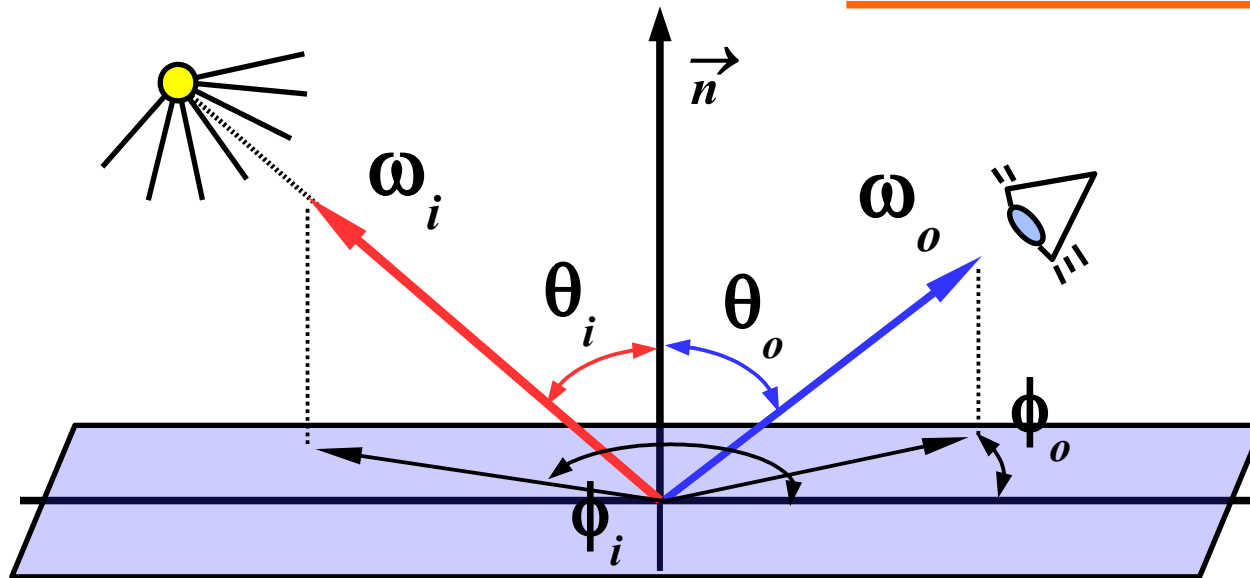
# Shading terms (components)



# BRDF formulation



BRDF function:  $\mathbf{R}^5 \rightarrow \mathbf{R}$   $f(\omega_i, \omega_o, \lambda)$



$$\underline{L_o(\omega_o)} = \int_{\Omega} \underline{f(\omega_i, \omega_o)} \cdot \underline{L_i(\omega_i)} (n \cdot \omega_i) \underline{d\omega_i}$$

$\swarrow$   
 $\cos \theta_i$



# BRDF plausibility

■ non-negative

$$f(\omega_i, \omega_o, \lambda) \geq 0$$

■ reciprocal

$$f(\omega_i, \omega_o, \lambda) = f(\omega_o, \omega_i, \lambda)$$

■ energy-conserving

$$\int_{\Omega} f(\omega_i, \omega_o) (n \cdot \omega_i) d\omega_i \leq 1$$

# History (**physical**/**empirical**)



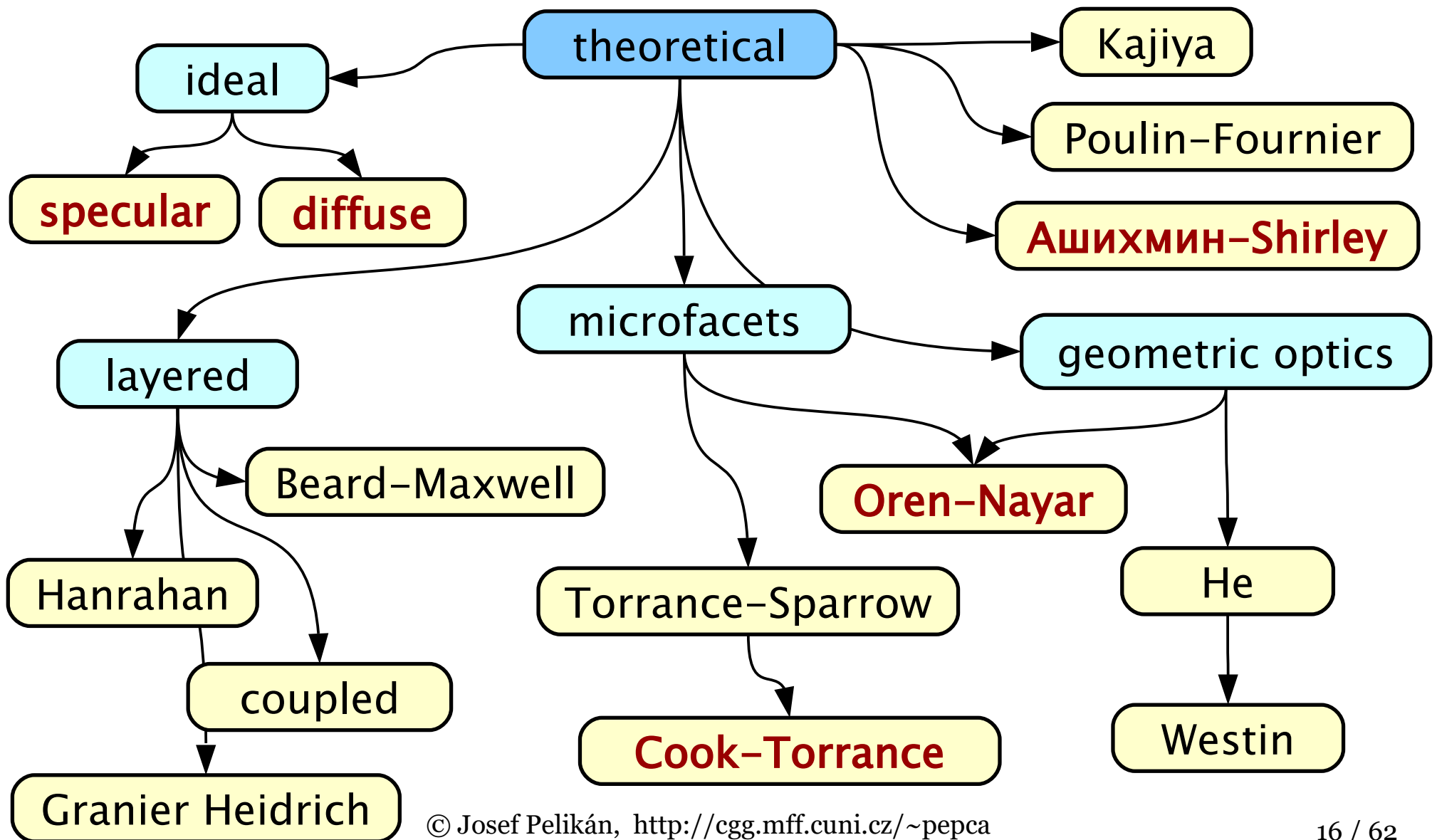
- **Beckmann, Spizzichino** (1963): electromagnetic wave reflection on rough surfaces (optics)
- **Torrance, Sparrow** (1967): off-specular reflections on rough surfaces (optics)
- **Phong** (1975): famous empirical model, used many decades
- **Blinn** (1977): first light reflection presentation at SIGGRAPH
- **Cook, Torrance** (1981): generalization, implementation, first physically based BRDF model in computer graphics



# History (**physical**/**empirical**)

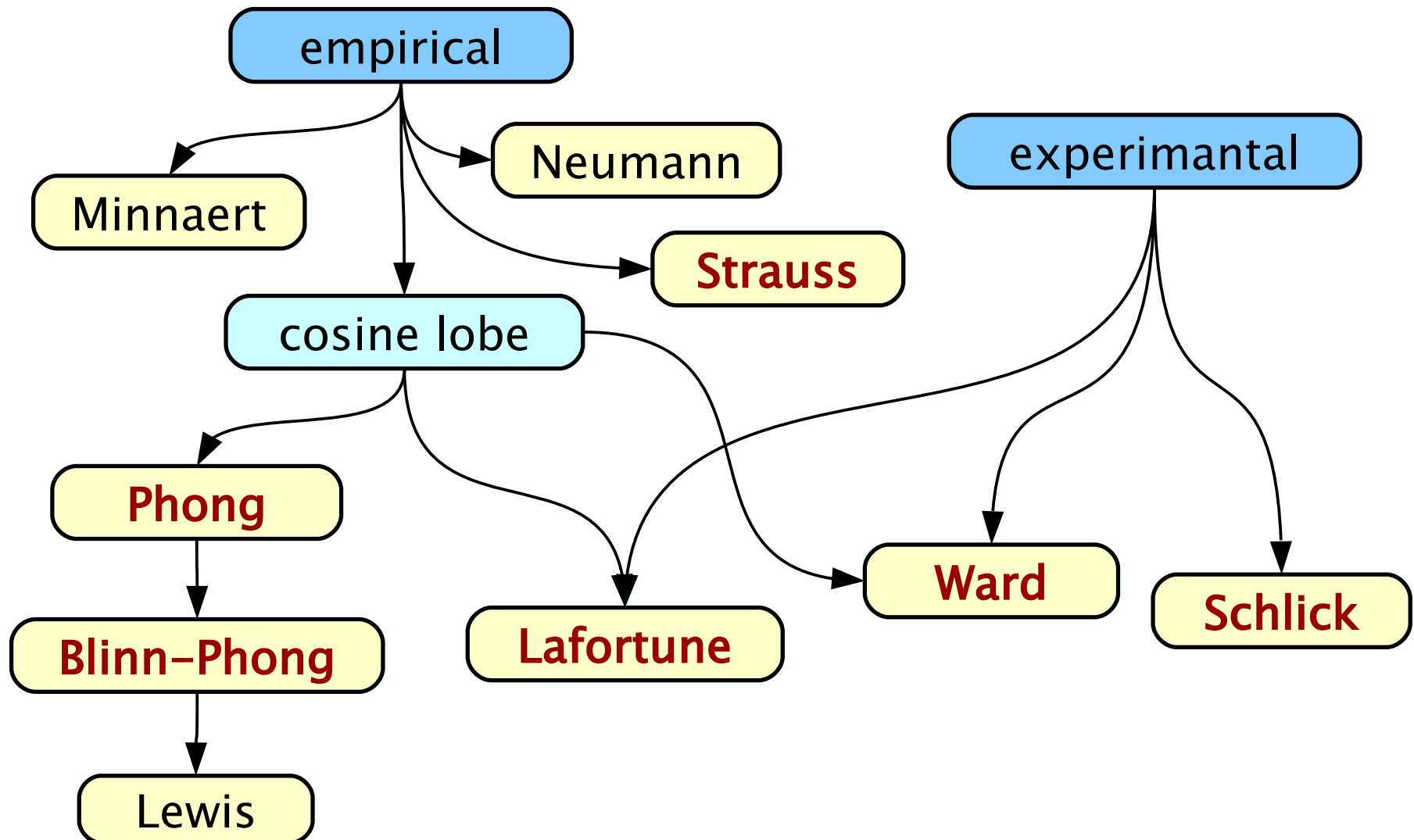
- **He** (1991): more complex wave optics, polarization, diffraction, interference..
- **Ward** (1992): anisotropic material, microfacets
- **Schlick** (1994): fast Fresnel formula approximation, two-layer reflectance model
- **Lafortune** (1997): multiple lobes, fitted to lab data
- **Ашихмин, Shirley** (2000): anisotropic Phong
- **Walter** (2007): microfacet refraction model (BSDF = Bidirectional Scattering Distribution Function)
- **Ашихмин, Bagher** (2007, 2012): models based on arbitrary microfacet distribution (measured..)

# BRDF ZOO I



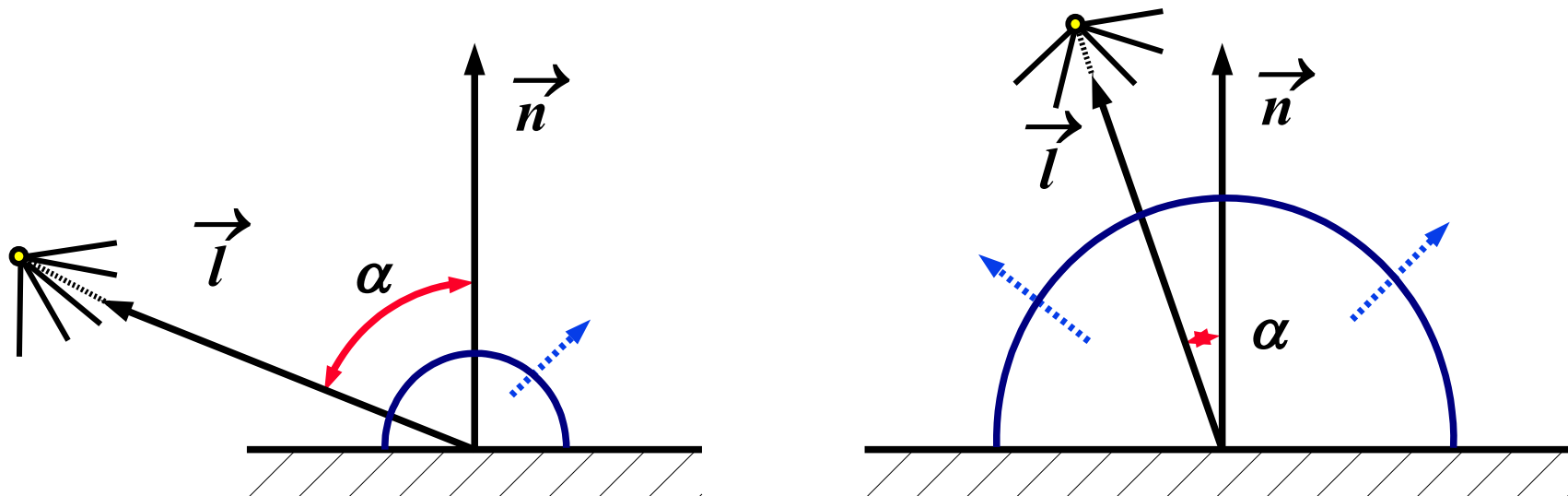


# BRDF ZOO II





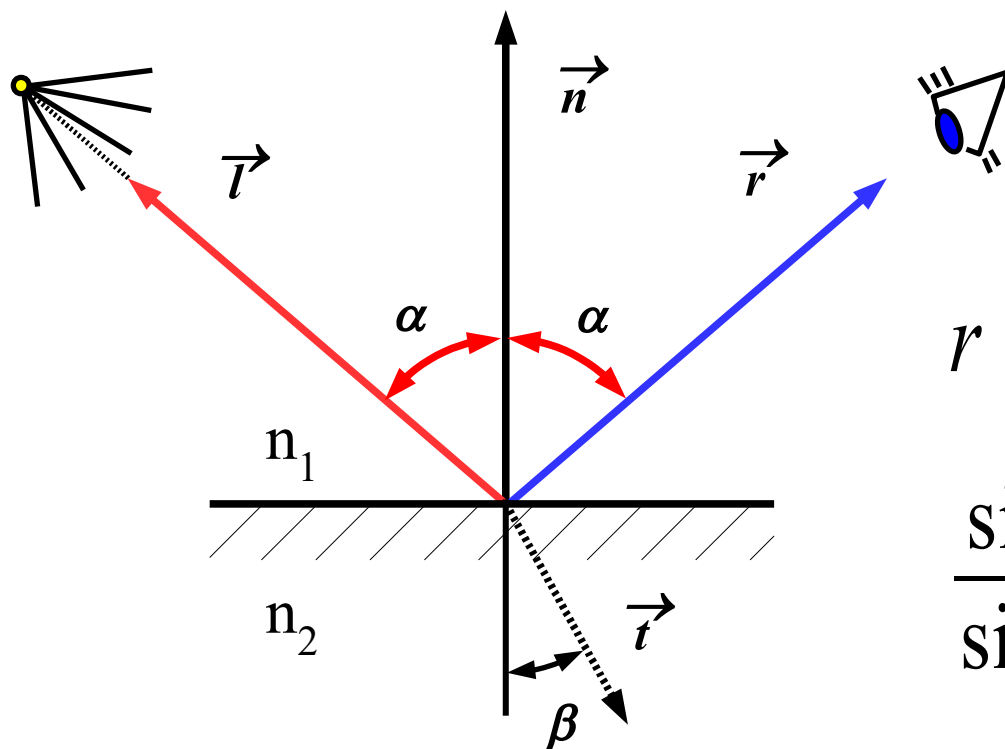
# Ideal diffusion



- **ideal diffuse material (Lambertian surface)**
  - ◆ reflection probability is constant
  - ◆ examples: furry surface, noisy microstructure w/o any pattern
- **Lambert law:** reflected intensity depends solely on  **$\cos \alpha$**



# Ideal (mirror) reflection

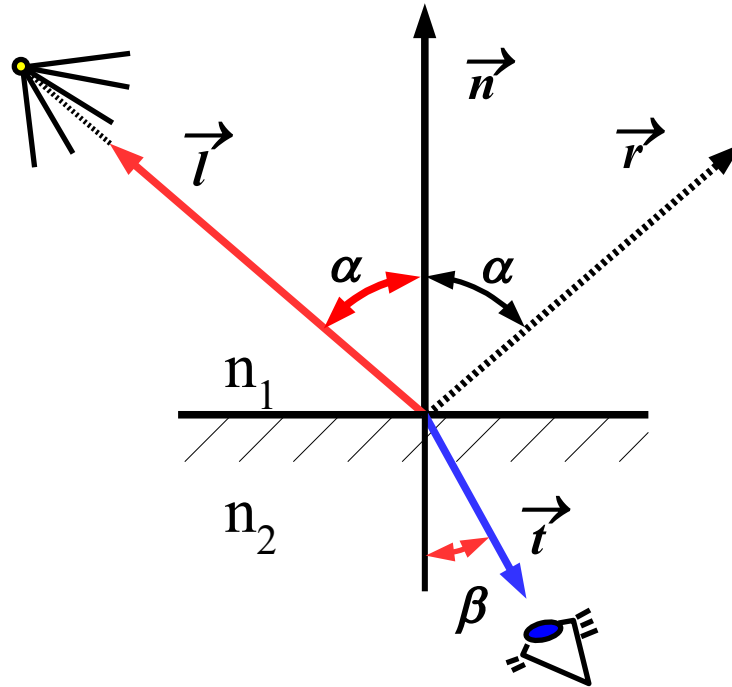


$$r = 2n(n \cdot l) - l$$

$$\frac{\sin \beta}{\sin \alpha} = \frac{n_1}{n_2} = n_{12}$$

- ratio of reflected and refracted light is determined by the **Fresnel equations** (19. century)

# Refraction (Snell's law, Ibn Sahl, 984)



$$\frac{\sin \beta}{\sin \alpha} = \frac{n_1}{n_2} = n_{12}$$

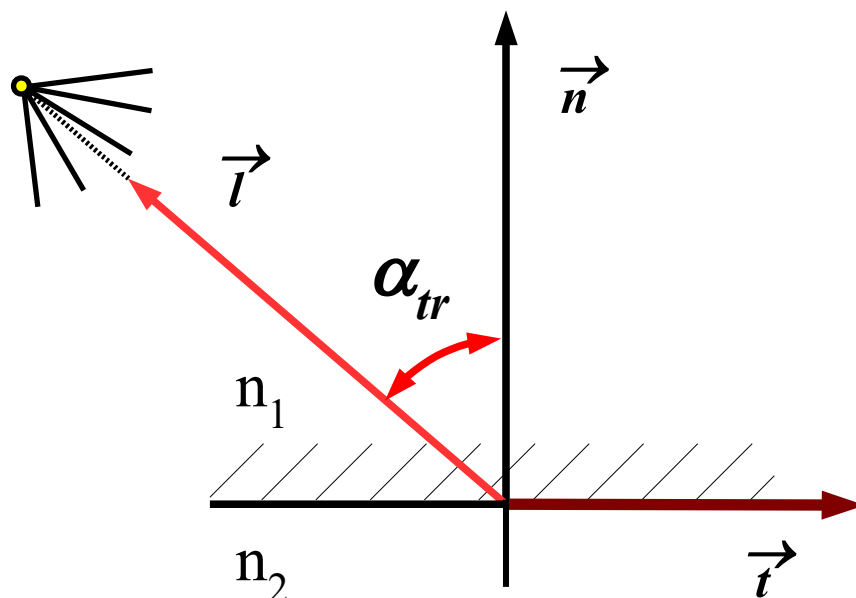
$$\cos \beta = \sqrt{1 - n_{12}^2 \sin^2 \alpha} = \sqrt{1 - n_{12}^2 \cdot (1 - (n \cdot l)^2)}$$

$$t = \left[ n_{12}(n \cdot l) - \sqrt{1 - n_{12}^2 \cdot (1 - (n \cdot l)^2)} \right] \cdot n - n_{12} \cdot l$$



# Total internal reflection

- going from more dense environment to less dense one ( $n_1 > n_2$ )
- for incident angles greater than **critical angle**  $\alpha_{tr}$  there is no refraction at all!



$$\sin \alpha_{tr} = \frac{n_2}{n_1}$$

# Fresnel equations (polarization)



- two **polarizations** (electric field perpendicular „**s**“ /senkrecht/ or parallel „**p**“ to the incident plane)
- reflectance “**R**“ and transmittance “**T**“ (power ratios):

$$R_s = \left[ \frac{\sin(\beta - \alpha)}{\sin(\beta + \alpha)} \right]^2 \quad T_s = 1 - R_s$$

$$R_p = \left[ \frac{\tan(\beta - \alpha)}{\tan(\beta + \alpha)} \right]^2 \quad T_p = 1 - R_p$$



# Fresnel equations (alternative)

- no need to compute angles (cosines are easy):

$$R_s = \left[ \frac{n_1 \cos \alpha - n_2 \cos \beta}{n_1 \cos \alpha + n_2 \cos \beta} \right]^2$$

$$R_p = \left[ \frac{n_1 \cos \beta - n_2 \cos \alpha}{n_1 \cos \beta + n_2 \cos \alpha} \right]^2$$



# Unpolarized light

- averaging values  $R_s$  a  $R_p$ :

$$R = \frac{1}{2} \frac{(a-u)^2 + b^2}{(a+u)^2 + b^2} \left[ \frac{(a+u-1/u)^2 + b^2}{(a-u+1/u)^2 + b^2} + 1 \right]$$

$$a^2 = \frac{1}{2} \left( \sqrt{(n_\lambda^2 - k_\lambda^2 + u^2 - 1)^2 + 4n_\lambda^2 k_\lambda^2} + n_\lambda^2 - k_\lambda^2 + u^2 - 1 \right)$$

$$b^2 = \frac{1}{2} \left( \sqrt{(n_\lambda^2 - k_\lambda^2 + u^2 - 1)^2 + 4n_\lambda^2 k_\lambda^2} - n_\lambda^2 + k_\lambda^2 - u^2 + 1 \right)$$

$$u = \cos \alpha = n \cdot l \quad n = n_\lambda - i k_\lambda \quad (\text{for dielectric } k_\lambda = 0)$$





# Dielectric (insulator) materials

■  $\mathbf{k}_\lambda = \mathbf{0} \Rightarrow$

$$a^2 = n_\lambda^2 + u^2 - 1 \quad b = 0$$

$$R = \frac{1}{2} \frac{(a-u)^2}{(a+u)^2} \left( \frac{[u(a+u)-1]^2}{[u(a-u)+1]^2} + 1 \right)$$



# Remarks (Fresnel)

- if  $\alpha = \pi/2$  (i.e.  $\mathbf{u} = \mathbf{o}$ ), then reflectance  $\mathbf{R}_\lambda(\mathbf{90}) = \mathbf{1}$  regardless of the wavelength  $\lambda$
- for perpendicular ray ( $\alpha = \mathbf{o}$ ):

$$R_0 = R_s = R_p = \left( \frac{n_2 - n_1}{n_2 + n_1} \right)^2$$

$$T_0 = T_s = T_p = 1 - R_0 = \frac{4n_1n_2}{(n_2 + n_1)^2}$$



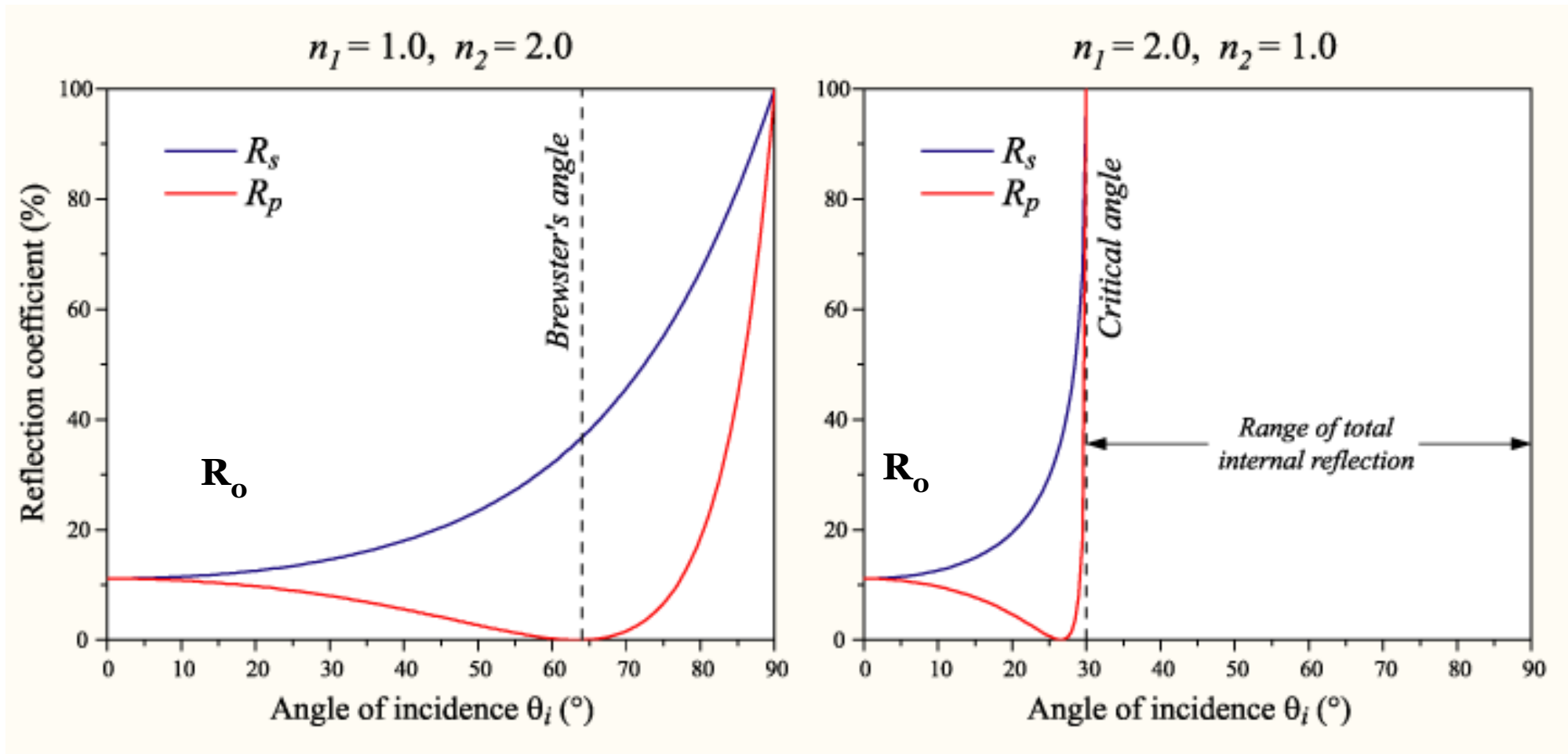
# Wavelength $\lambda$

For  $\mathbf{l}$  and  $\mathbf{v}$  perpendicular to the surface (i.e.  $\alpha = \mathbf{0}$ ):

$$\mathbf{F}(\lambda, \mathbf{0}) = \left( \frac{\mathbf{n}_\lambda - 1}{\mathbf{n}_\lambda + 1} \right)^2 \quad \mathbf{a} \quad \mathbf{n}_\lambda = \frac{1 + \sqrt{\mathbf{F}(\lambda, \mathbf{0})}}{1 - \sqrt{\mathbf{F}(\lambda, \mathbf{0})}}$$

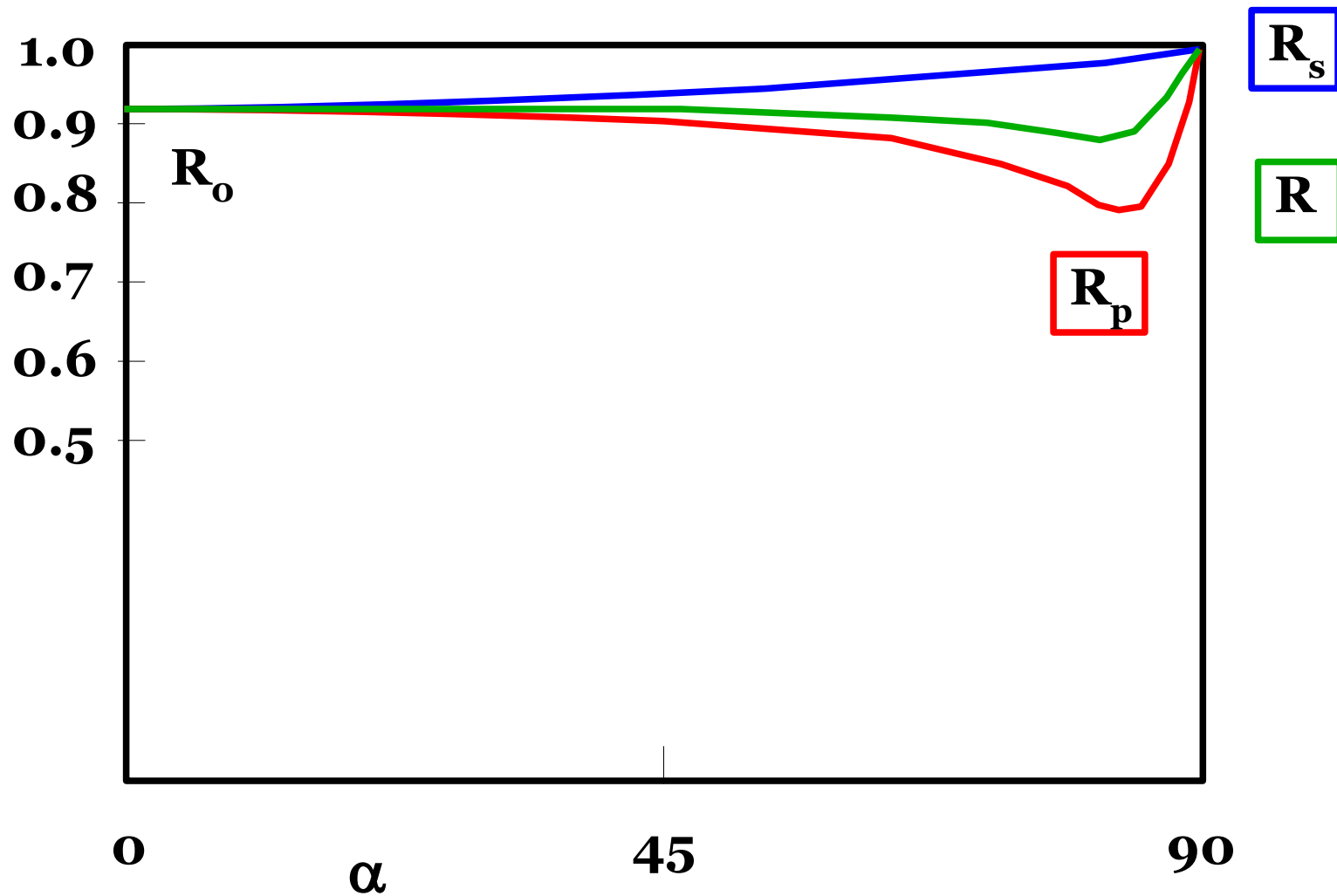
- quantities  $\mathbf{F}_\lambda(\mathbf{0})$  were measured in labs for many real materials (both conductors and insulators)
  - so we know the  $\mathbf{n}_\lambda$  indices
- specular reflection **depends on  $\lambda$**  (except for  $\alpha = \pi/2$  )

# Reflectance – dielectric material



(cc) Ulflund, Wiki

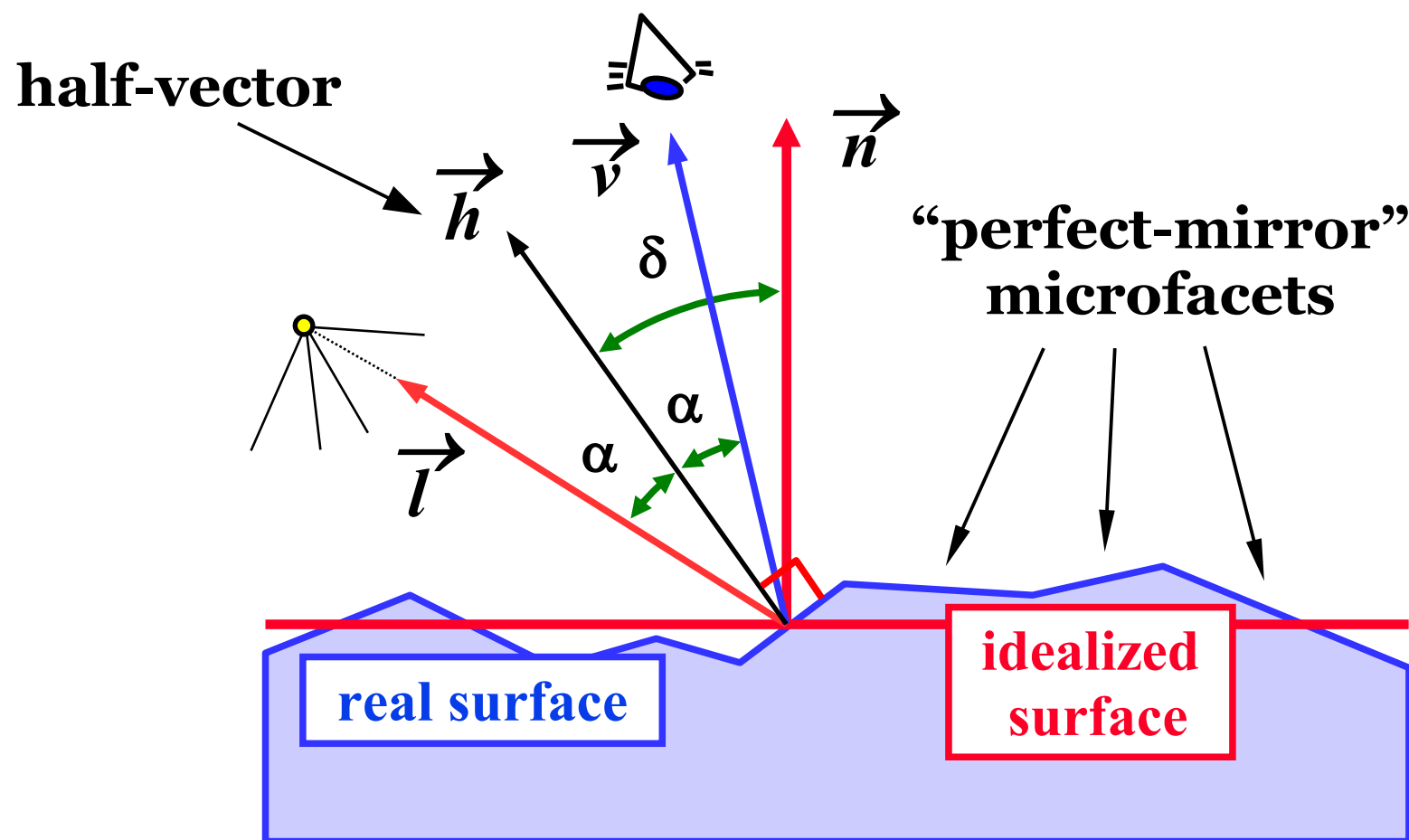
# Reflectance – metal





# Microfacet theory

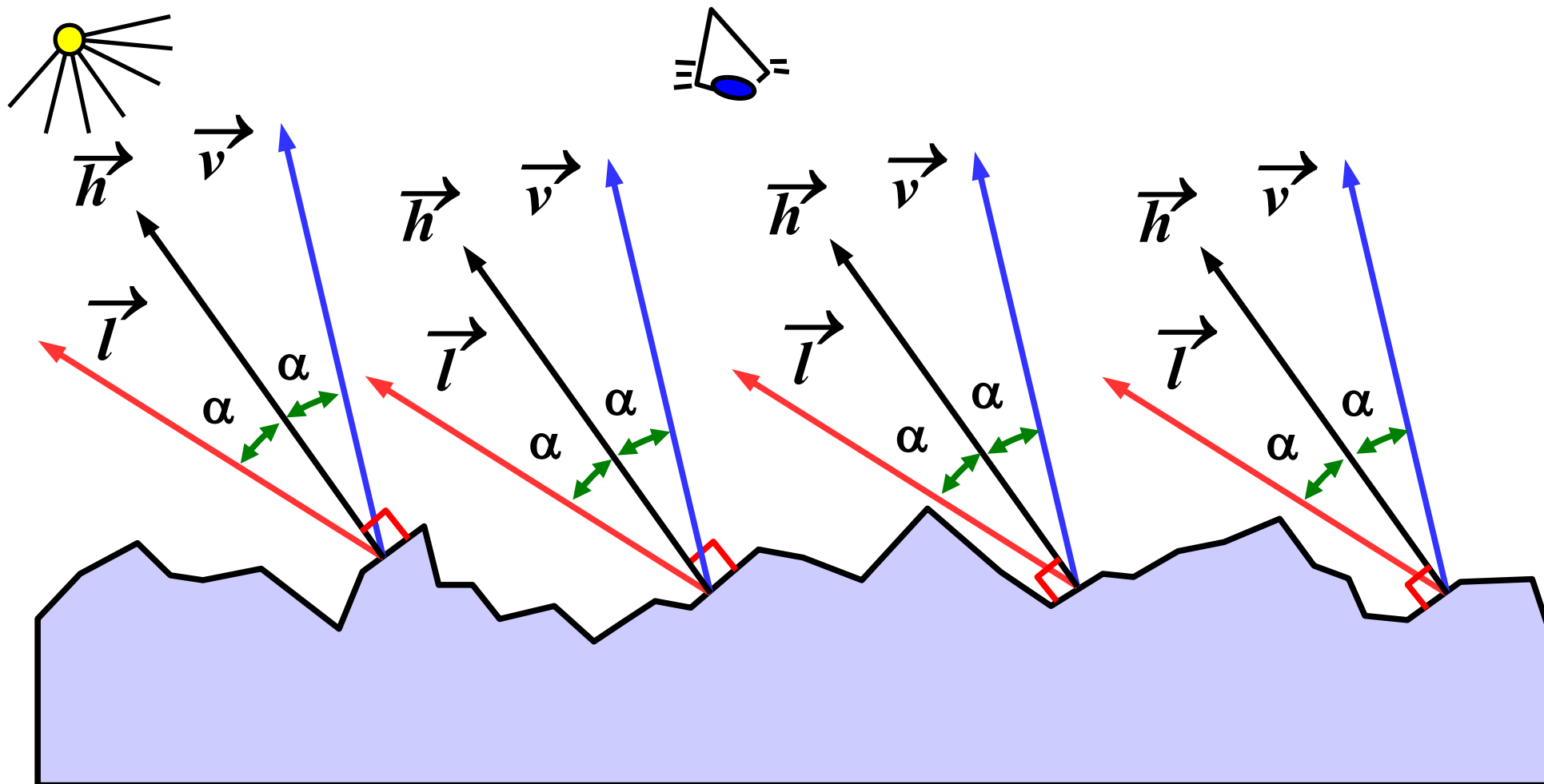
- Beckmann, Spizzichino (63), Torrance, Sparrow (67)



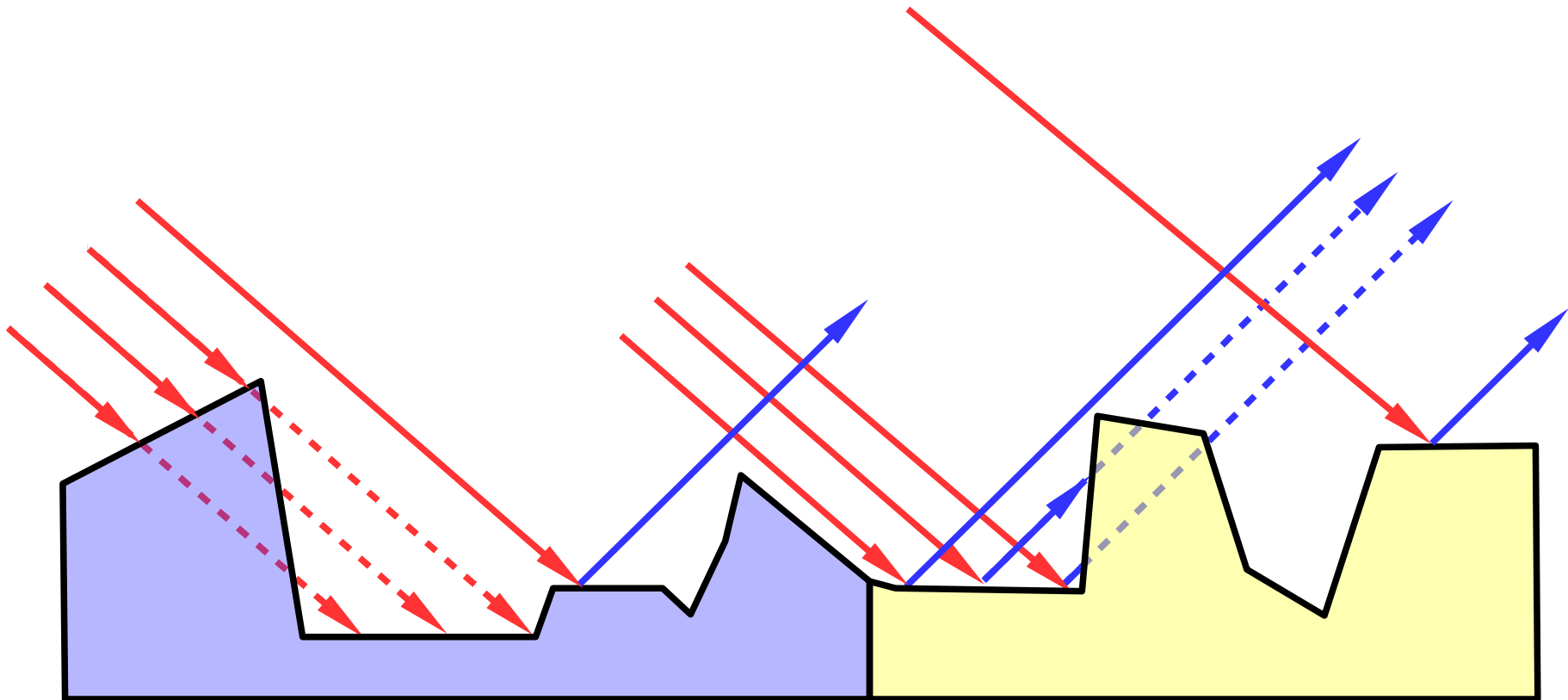
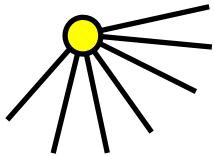


# Perfect reflection for half-angle

Only ideal **half-angle microfacets** can contribute !



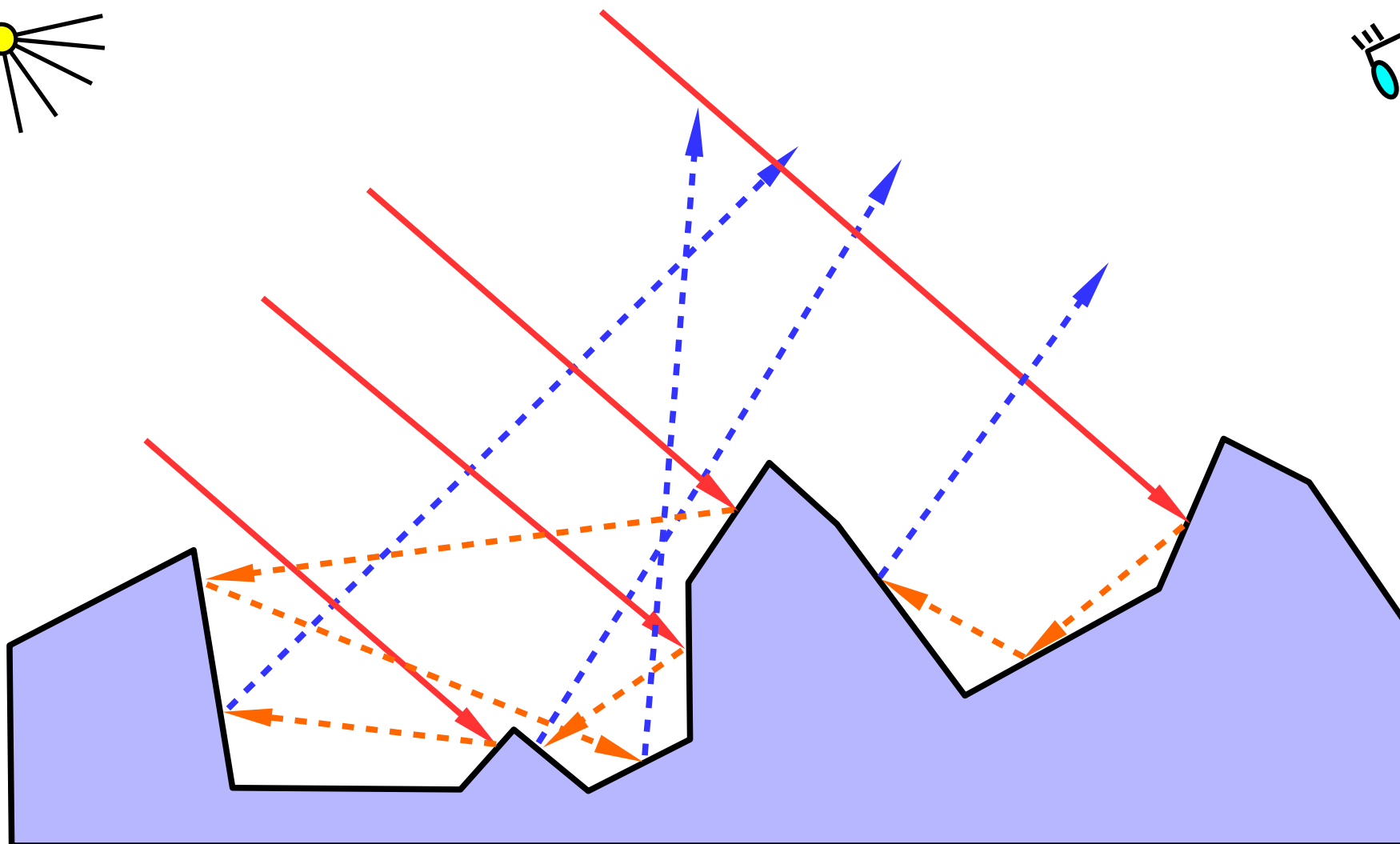
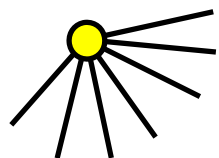
# Shadowing and masking







# Multiple bounces are lost





# Microfacet specular BRDF

$$R_{\lambda}(h) = \frac{F_{\lambda}(\alpha)}{4} \cdot \frac{D(h) \cdot G(l, v, h)}{(n \cdot l)(n \cdot v)}$$

- $R_{\lambda}(h)$  ... specular reflectance for wavelength  $\lambda$
- $F_{\lambda}(\alpha)$  ... Fresnel ideal reflectance for wavelength  $\lambda$  and incident angle  $\alpha$
- $D(h)$  ... microfacet PDF (“how many microfacets” have  $h$  as a normal vector)
- $G(l, v, h)$  ... geometric factor (shadowing & masking)



# Fresnel term F

Fresnel term for unpolarized light

$$F(\lambda, \beta) = \frac{1}{2} \cdot \frac{(\mathbf{g} - \mathbf{c})^2}{(\mathbf{g} + \mathbf{c})^2} \left\{ 1 + \frac{[\mathbf{c}(\mathbf{g} + \mathbf{c}) - \mathbf{1}]^2}{[\mathbf{c}(\mathbf{g} - \mathbf{c}) - \mathbf{1}]^2} \right\}$$

for  $\mathbf{c} = \cos \beta = (\mathbf{V} \cdot \mathbf{H})$ ,

$$\mathbf{g}^2 = \mathbf{n}_\lambda^2 + \mathbf{c}^2 - 1$$

- $\mathbf{n}_\lambda$  ... index of refraction for wavelength  $\lambda$
- ◆ for conductors  $\mathbf{n}_\lambda' = \mathbf{n}_\lambda - \mathbf{i} \kappa_\lambda$  ( $\kappa_\lambda$  ... absorption coeff.)



# Fresnel – base specular color

Metal	F(0) [linear]	F(0) [sRGB]
Titanium	0.542, 0.497, 0.449	194, 187, 179
Chromium	0.549, 0.556, 0.554	196, 197, 196
Iron	0.562, 0.565, 0.578	198, 198, 200
Nickel	0.660, 0.609, 0.526	212, 205, 192
Platinum	0.673, 0.637, 0.585	214, 209, 201
Copper	0.955, 0.638, 0.538	250, 209, 194
Palladium	0.733, 0.697, 0.652	222, 217, 211
Zinc	0.664, 0.824, 0.850	213, 234, 237
Gold	1.022, 0.782, 0.344	255, 229, 158
Aluminum	0.913, 0.922, 0.924	245, 246, 246
Silver	0.972, 0.960, 0.915	252, 250, 245





# Schlick's approximation

- Fresnel term for other angles, based on  $F_\lambda(\mathbf{o}) = \mathbf{c}$

$$F_{\text{schlick}}(\mathbf{c}, \mathbf{l}, \mathbf{h}) = \mathbf{c} + (1 - \mathbf{c}) (1 - (\mathbf{l} \cdot \mathbf{h}))^5$$



# Any angle, any $\lambda$ (R. Hall)

- let's assume we have a **base material color**  $F_\lambda(\mathbf{o})$  and an **angle-function** for some (standard)  $\lambda_0$ 
  - ◆ set of wavelengths can be limited (3÷6 components)

$$F_\lambda(\alpha) \approx F_\lambda(0) + (1 - F_\lambda(0)) \frac{\max(0, F_{\lambda_0}(\alpha) - F_{\lambda_0}(0))}{1 - F_{\lambda_0}(0)}$$



# Normal distribution function

Fast and simple formula – Gaussian distribution:

$$\underline{D(h, m)} = \chi(n \cdot h) (n \cdot h) \cdot e^{-\left(\frac{\delta}{m}\right)^2}$$

- $\chi(a) = a > 0 ? 1 : 0$
- $\cos \delta = \mathbf{n} \cdot \mathbf{h}$
- $m$  ... “surface roughness” (standard deviation of the surface slope)
  - $< 0.1$  ... very smooth
  - $> 0.8$  ... rough (almost diffuse)

# Beckmann's distribution (~normalised)

$$\begin{aligned} \underline{D_{be}(h, m)} &= \frac{\chi(n \cdot h)}{\pi m^2 (n \cdot h)^4} e^{-\left(\frac{\tan \delta}{m}\right)^2} \\ &= \frac{\chi(n \cdot h)}{\pi m^2 (n \cdot h)^4} e^{\frac{(n \cdot h)^2 - 1}{m^2 (n \cdot h)^2}} \end{aligned}$$





# Blinn–Phong (normalised)

$$\underline{D_{bp}(h, m)} = \chi(n \cdot h) \frac{m+2}{2\pi} (n \cdot h)^m$$

# Trowbridge & Reitz



$$\underline{D_{tr}(h, m)} = \frac{\chi(n \cdot h) m^2}{\pi \left( (n \cdot h)^2 (m^2 - 1) + 1 \right)^2}$$

- $m$  can be greater than 1

# GGX (Walter et al. 2007)



$$\underline{D_{GGX}}(h, m) = \frac{\chi(n \cdot h) m^2}{\pi (n \cdot h)^4 (m^2 + \tan^2 \delta)^2}$$



# Isotropic Ward (1992)

$$\underline{D_{wiso}(h, m)} = \frac{\chi(n \cdot h)}{\pi m^2} e^{-\frac{\tan^2 \delta}{m^2}}$$



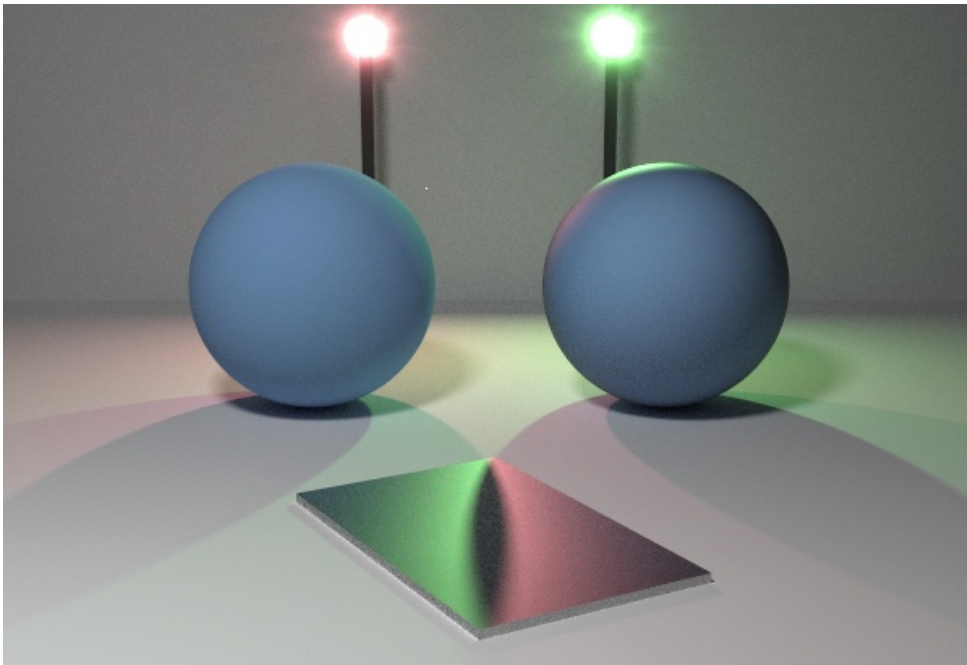
# Anisotropic Ward (1992)

$$\underline{D_w(h, m_x, m_y)} = \frac{\chi(n \cdot h)}{\pi m_x m_y} e^{-\tan^2 \delta \left( \frac{\cos^2 \phi_h}{m_x^2} + \frac{\sin^2 \phi_h}{m_y^2} \right)}$$

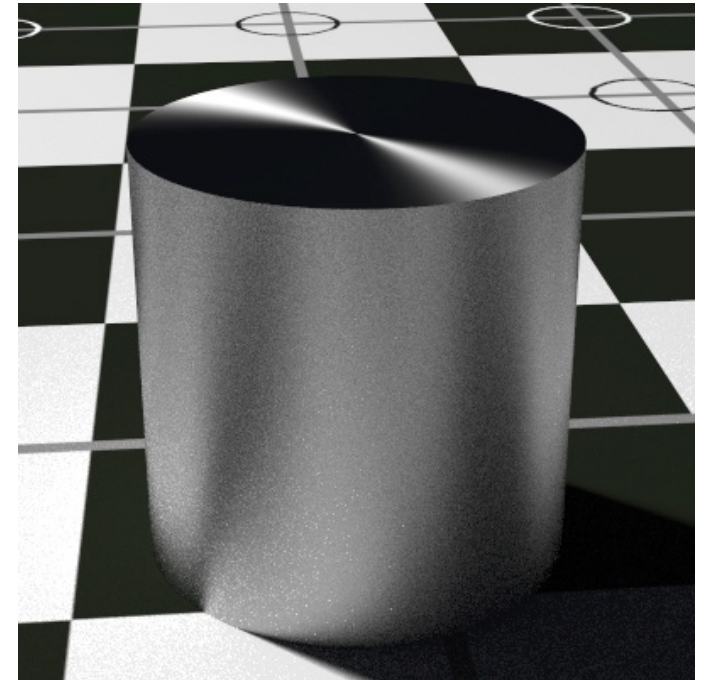
- $\phi_h$  ... azimuth angle of the half-vector

# Ашихмин–Shirley anisotropic (2000)

$$\underline{D_{as}(h, e_x, e_y)} = \sqrt{(e_x + 1)(e_y + 1)} (h \cdot n)^{e_x \cos^2 \phi_h + e_y \sin^2 \phi_h}$$



(c) Ашихмин, 2000





# Material blends

Idea of **blending several materials together**  
makes sense:  $\mathbf{m}_1 \dots \mathbf{m}_k$

$$\mathbf{D}(\alpha) = \sum_{i=1}^k \mathbf{w}_i \cdot \mathbf{D}(\mathbf{m}_i, \alpha)$$

- $\mathbf{w}_i$  ... weight coefficients

$$\sum \mathbf{w}_i = \mathbf{1}$$

# Geometric term G (Cook–Torrance)

Compensation for masking and shadowing

$$\underline{G_{ct}}(l, v, h) = \min\left(1, \frac{2(n \cdot h)(n \cdot v)}{(v \cdot h)}, \frac{2(n \cdot h)(n \cdot l)}{(v \cdot h)}\right)$$

Cook and Torrance assumed infinitely long  
“**V**”-shaped grooves (not exactly true)

- optimized: **Kelemen**, more accurate: **Smith**





# G alternative (GGX)

$$\underline{G_{GGX}}(l, v, h) = \chi\left(\frac{v \cdot h}{v \cdot n}\right) \frac{2}{1 + \sqrt{1 + m^2 \tan^2 \theta_v}}$$

- **$m$**  ... roughness from the GGX distribution

# Cheap G (Kelemen–Szirmay–Kalos)

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$$\frac{G_{KSK}(l, v, h)}{(n \cdot l)(n \cdot v)} \approx \frac{1}{(l \cdot h)^2}$$

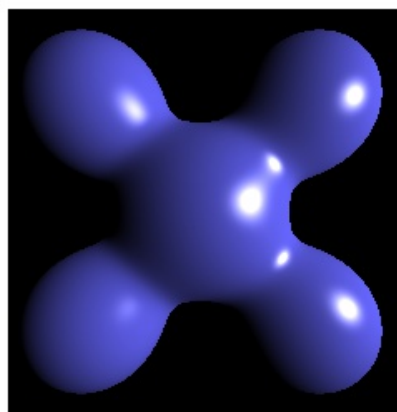


# Blinn's contribution

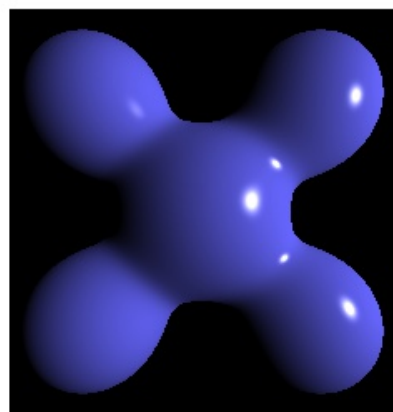
- Blinn-Phong model (here  $\beta = \angle v, r$ )
  - ◆ light source and viewer in infinity:

$$\cos^h \beta \approx \cos^{4h} \beta / 2$$

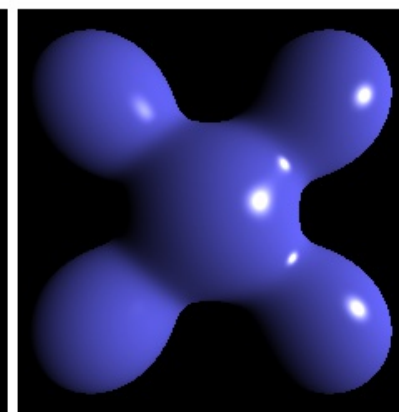
$$(R_i \cdot V)^h \approx (H_i \cdot N)^{4h}$$



**Blinn-Phong**



**Phong**



**Blinn-Phong**  
(higher exponent)

(cc) Wiki



# Schlick's contribution

- **Christophe Schlick** (1994) was experimenting with approximate substitution in Fresnel term
- fraction instead of power function

$$S^h \approx S / (h - hS + S)$$

- ◆ slightly less sharp highlight compared to Blinn-Phong

- Fresnel term substitute (32× faster, error <1%)

$$R_{\text{schlick}}(c, l, n) = c + (1 - c) (1 - (l \cdot n))^5$$



# Lafortune model (1997)

## ■ Generalized cosine lobe model

- ◆ derived using Householder matrix ( $3 \times 3$ )

1. classical specular term (Phong) ..

$$f_r(l, v) = \rho_s C_s \cos^h \beta$$

2. .. rewritten using Householder matrix notation

$$\begin{aligned} f_r(l, v) &= \rho_s C_s (r \cdot v)^h \\ &= \rho_s C_s \left[ l^T (2n n^T - I) v \right]^h \end{aligned}$$



# General plausible cosine lobe

- Householder matrix  $\mathbf{M}$  ( $3 \times 3$ )
  - for reciprocity it must be **symmetrical**

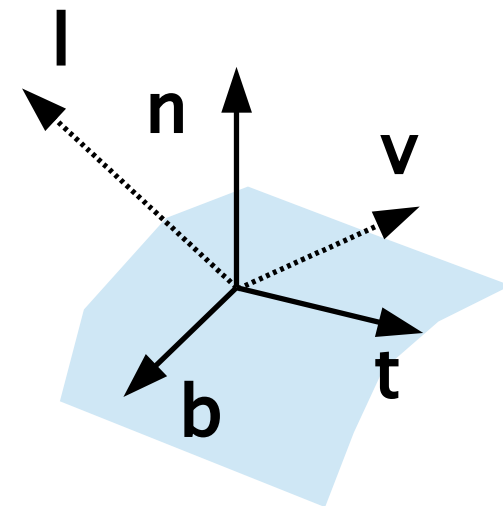
$$f_r(l, v) = \rho_s \left[ l^T M v \right]^h$$

- SVD decomposition of matrix  $\mathbf{M}$ :

$$f_r(l, v) = \rho_s \left[ l^T Q^T D Q v \right]^h$$

- $\mathbf{Q}$  ... coordinate transform,  $\mathbf{D}$  ... diagonal matrix

$$f_r(l, v) = \rho_s \left[ C_b l_b v_b + C_t l_t v_t + C_n l_n v_n \right]^h$$





# Cosine lobe options

- Phong lobe:  $-C_b = -C_t = C_n = \sqrt{C_s}$
- general isotropic reflection:  $C_b = C_t$
- anisotropy:  $C_b \neq C_t$
- isotropic diffuse term:  $C_b = C_t = \mathbf{0}, C_s = (\mathbf{h} + \mathbf{2}) / 2\pi$
- off-specular reflection:  $C_n < -C_b = -C_t$
- retro-reflections:  $C_b, C_t, C_n > \mathbf{0}$

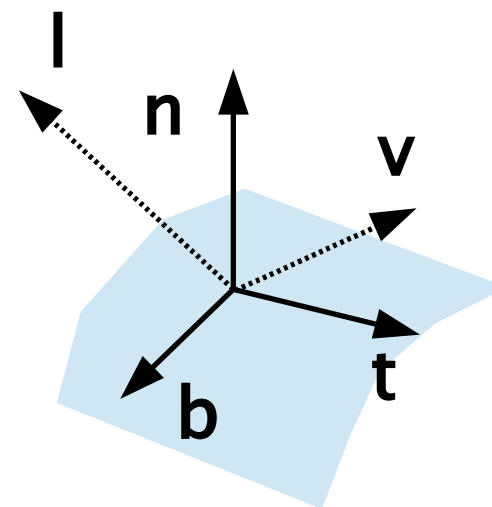


# Compound model

- Superposition of several lobes

- ◆ each one is defined by:  $C_{b,i}$ ,  $C_{t,i}$ ,  $C_{n,i}$ ,  $h_i$

$$f_r(l, v) = \sum_i \left[ C_{b,i} l_b v_b + C_{t,i} l_t v_t + C_{n,i} l_n v_n \right]^{h_i}$$

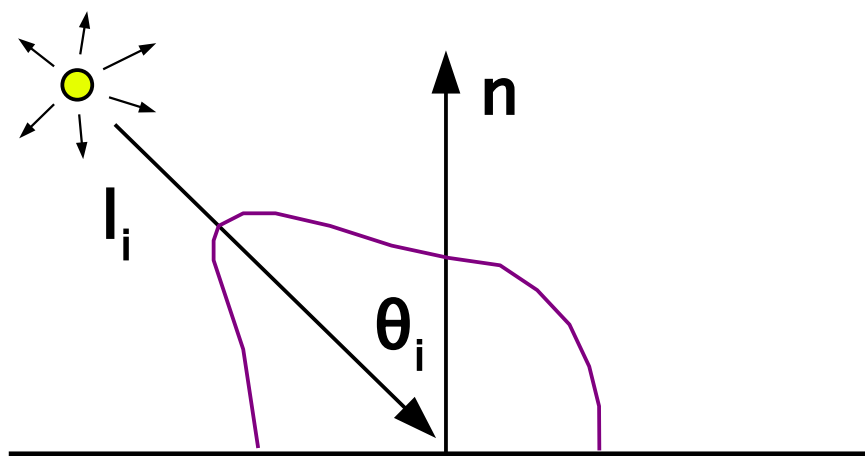






# Lambert law is not perfect..

- pure "cosine" surface is not as common in the nature
  - ◆ rough, grainy surfaces (sandpaper, sand, etc.)
  - ◆ **full moon** – contours should be darker but actually **they are not !**
  - ◆ "back-scattering" effect, reflecting (passive) taillights





# Oren–Nayar model

- based on **microfacet idea**
  - ◆ diffuse reflection on microfacets
  - ◆ simplified formulas – only most important ones

$$E_d = \frac{\rho}{\pi} \cdot E_0 \cdot \cos(\theta_i) \cdot (A + B \cdot \max(0, \cos(\phi_o - \phi_i))) \cdot \sin(\alpha) \cdot \tan(\beta)$$

$\theta_i$  incoming angle ( $\angle l_i, n$ )

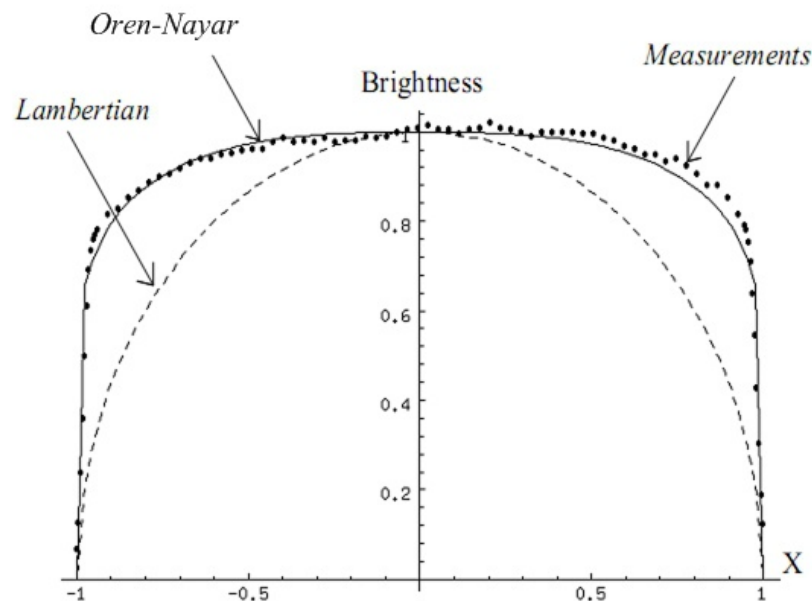
$\theta_o$  outgoing angle ( $\angle v, n$ )

$\Phi_i$  incoming azimuth of  $\omega_i$

$\Phi_o$  outgoing azimuth of  $\omega_o$

$\alpha$   $\max(\theta_i, \theta_o)$

$\beta$   $\min(\theta_i, \theta_o)$





# Oren–Nayar, final formula

$$E_d = (C_L \circ C_D) \cdot \cos(\theta_i) \cdot (A + B \cdot \max(0, \cos(\phi_o - \phi_i))) \cdot \sin(\alpha) \cdot \tan(\beta)$$

$$A = 1 - 0.5 \cdot \frac{\sigma^2}{\sigma^2 + 0.33} \quad (\text{value in denominator} - \text{up to } 0.57)$$

$$B = 0.45 \cdot \frac{\sigma^2}{\sigma^2 + 0.09}$$

$\sigma$     **roughness:** mean value of  $h$  (see Cook-Torrance)

$C_L$     light source color

$C_D$     material color



# Oren-Nayar - samples



Real Image

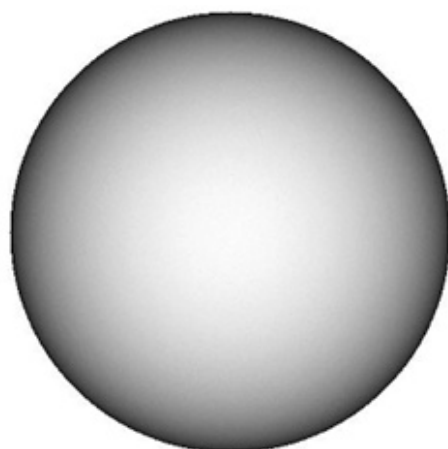


Lambertian Model

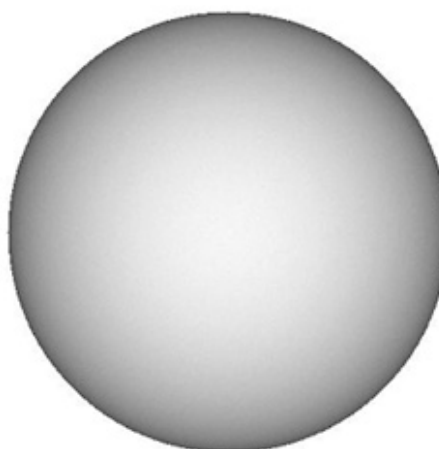


Oren-Nayar Model

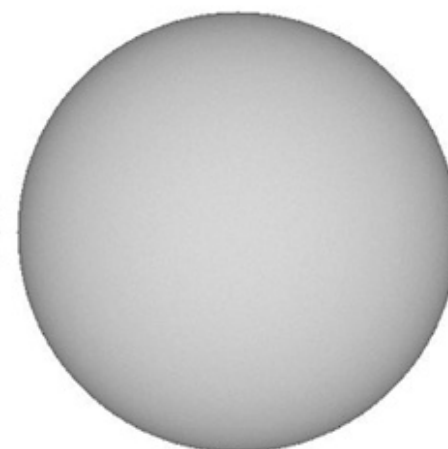
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$\sigma = 0$



$\sigma = 0.1$



$\sigma = 0.3$



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